

# Beam Loading in the HPRF Cavity

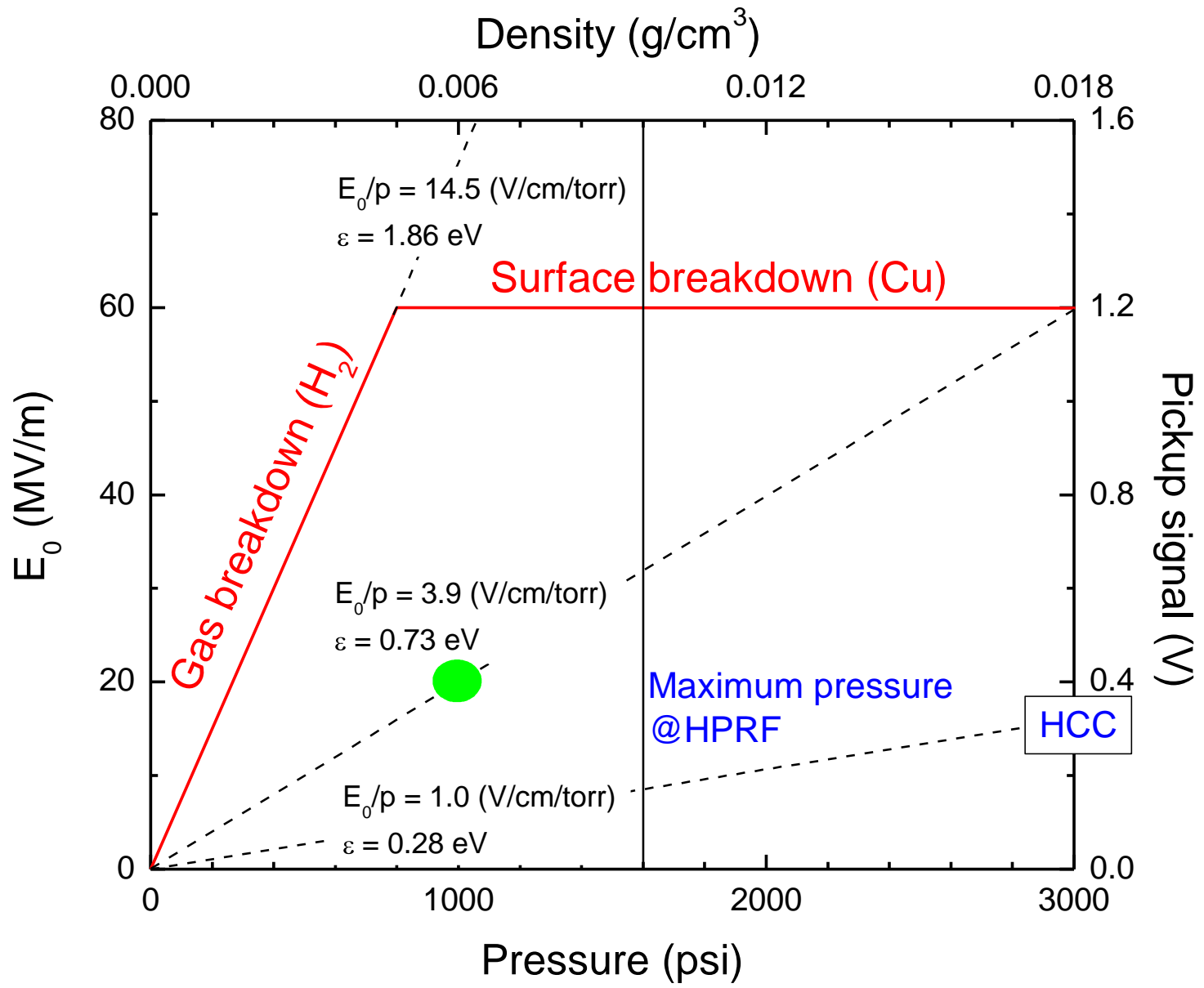
2009. 2. 18.

Moses Chung

APC, Fermilab

with

Alvin Tollestrup, Andreas Jansson, Katsuya Yonehara



## Without beam

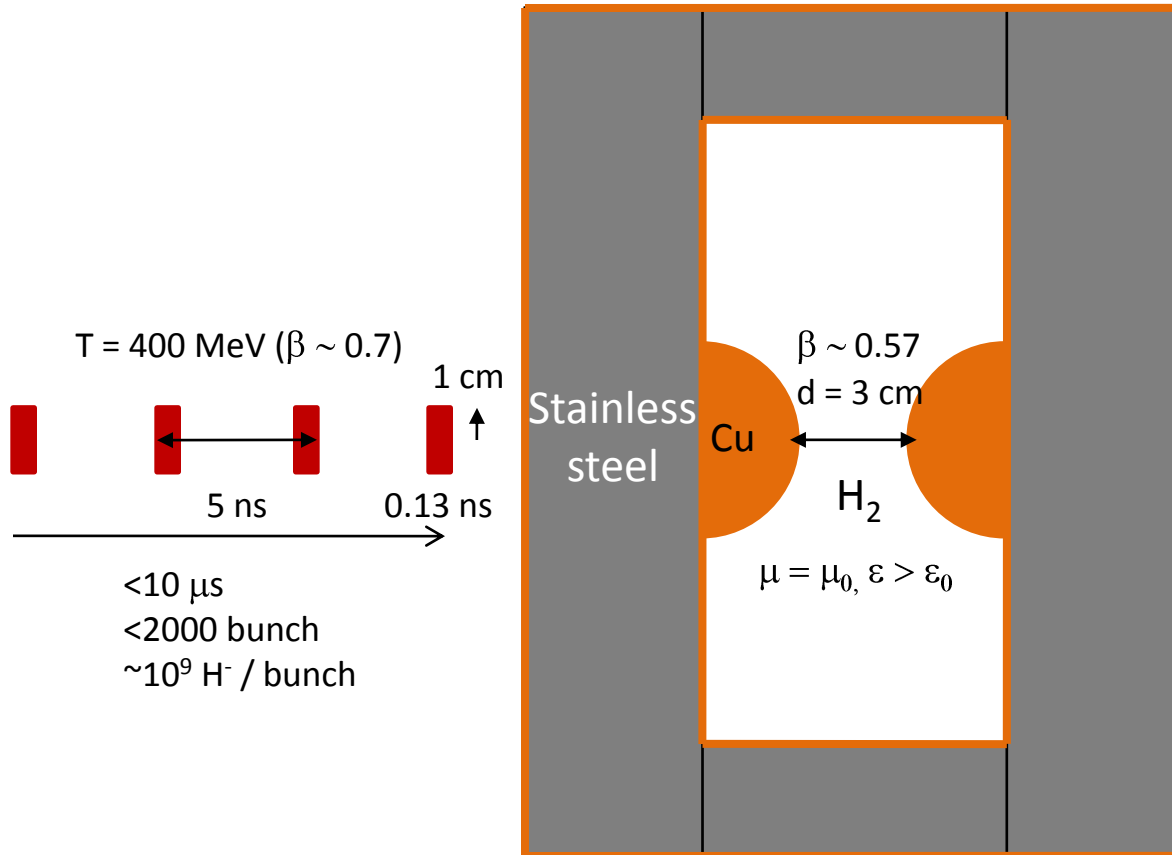
Surface-emitted  
electrons  
(K. Yonehara)

## Beam with Material

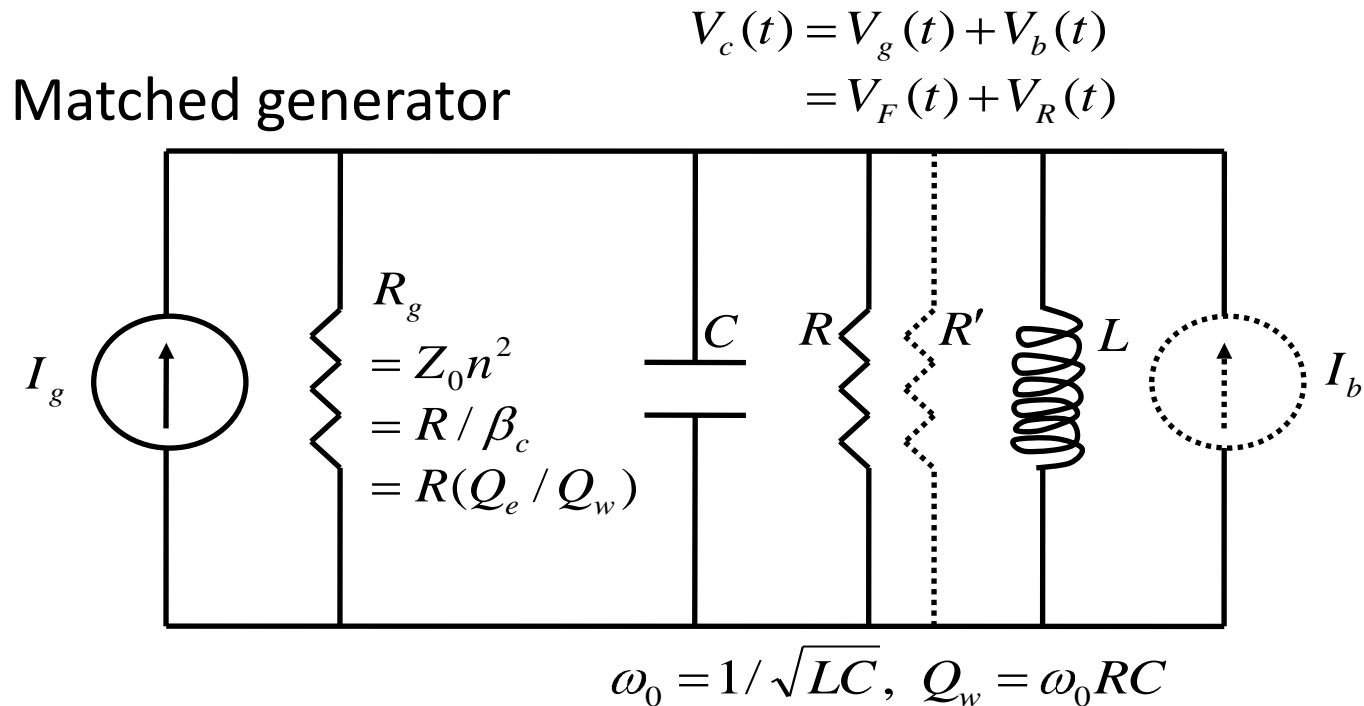
Secondary electrons  
 $\sim 2/1000$  proton  
(I. Rakhno)

## Beam in the Cavity

Beam-induced electrons  
 $\sim 1000/\text{proton}$   
(A. Tollestrup)



# Equivalent Circuit Model



$$\left\{ \frac{d^2}{dt^2} + \omega_0 \left( \frac{1}{Q_L} + \Delta \left( \frac{1}{Q} \right) \right) + \omega_0^2 \right\} V_c = 2 \frac{\omega_0}{Q_e} \frac{dV_F}{dt} - \frac{\omega_0}{2} \left[ \frac{R}{Q} \right] \frac{dI_b}{dt}$$

Additional damping term by  
beam-induced electrons (as plasma people do)

Additional driving term by  
beam itself (as RF people do)

# Model Equation

In the slowly-varying envelope approximation:  $|d\tilde{V}_c / dt| \ll |\omega \tilde{V}_c|$

$$\frac{d\tilde{V}_c}{d\tau} + (1 - j \tan \psi + \gamma) \tilde{V}_c = \frac{1}{2} Q_L \left[ \frac{R}{Q} \right] (\tilde{I}_g - \tilde{I}_b)$$

beam-induced electrons
beam itself

$$V_c(t) = \text{Re}(\tilde{V}_c(t)e^{j\omega t}), \quad |\tilde{V}_c| \approx E_0 T d, \quad T = \text{transit-time factor} \approx 0.96$$

$$\tau = \frac{t}{T_f}, \quad T_f = 2Q_L / \omega_0 = \text{filling time}$$

$$\tan \psi = Q_L \left( \frac{\omega_0}{\omega} - \frac{\omega}{\omega_0} \right)$$

$$\gamma = Q_L \Delta \left( \frac{1}{Q} \right)$$

$$\left[ \frac{R}{Q} \right] = \frac{|\tilde{V}_c|^2}{\omega_0 U} \approx 377 \times T^2 \frac{d}{R_w} \sim 129 \, \Omega$$

# 1. Beam Itself

# Beam Current

$$I_{DC} = \frac{q_b}{T_b} = \frac{Ne}{k(2\pi/\omega)} \approx \frac{10^9 (1.6 \times 10^{-19})}{4/(805 \text{ MHz})} \approx 32 \text{ mA}$$

For  $t_b$  (bunch length  $\approx 0.13 \text{ ns}$ )  $\ll T_b$  (bunch spacing  $\approx 5 \text{ ns}$ )

$$I_b(t) = q_b \sum_{m=-\infty}^{\infty} \delta(t - mT_b) = \frac{2q_b}{T_b} \text{Re} \left[ \frac{1}{2} + \sum_{n=1}^{\infty} e^{j\omega t \left( \frac{n}{k} \right)} \right]$$

For a given harmonics

$$\tilde{I}_b = \frac{2q_b}{T_b} = 2I_{DC}$$

Beam-induced voltage **in steady state with  $\gamma=0$**

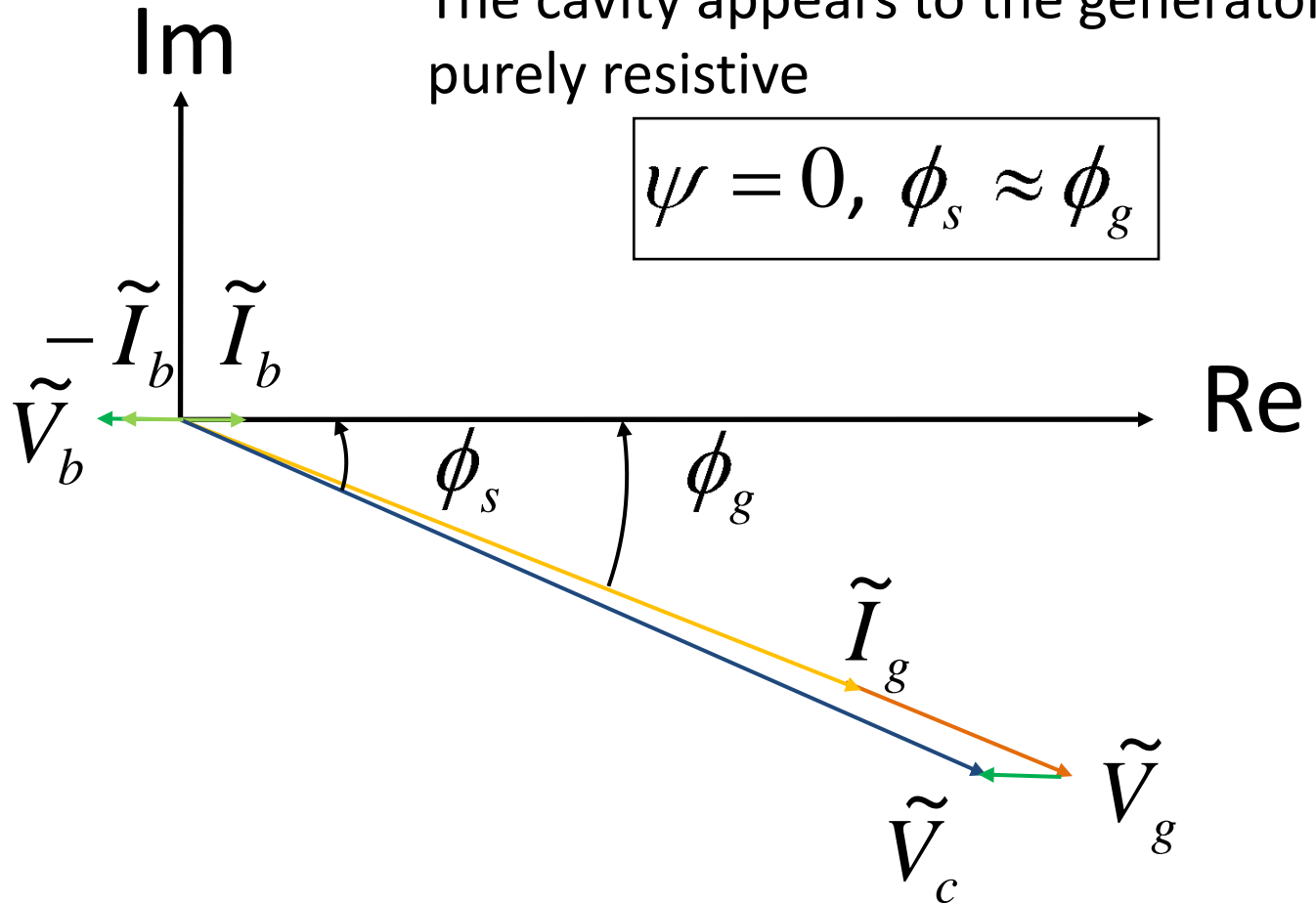
$$\tilde{V}_b = -\cos\psi e^{j\psi} \frac{1}{2} Q_L \left[ \frac{R}{Q} \right] \tilde{I}_b = -Z_{||}(\omega) \tilde{I}_b$$

$$\left| \tilde{V}_b \right| \sim 25 \text{ kV} \ll \left| \tilde{V}_c \right| \approx 20 \text{ MV/m} \times 3 \text{ cm} \times 0.96 \approx 580 \text{ kV}$$

# If $V_b$ negligible...

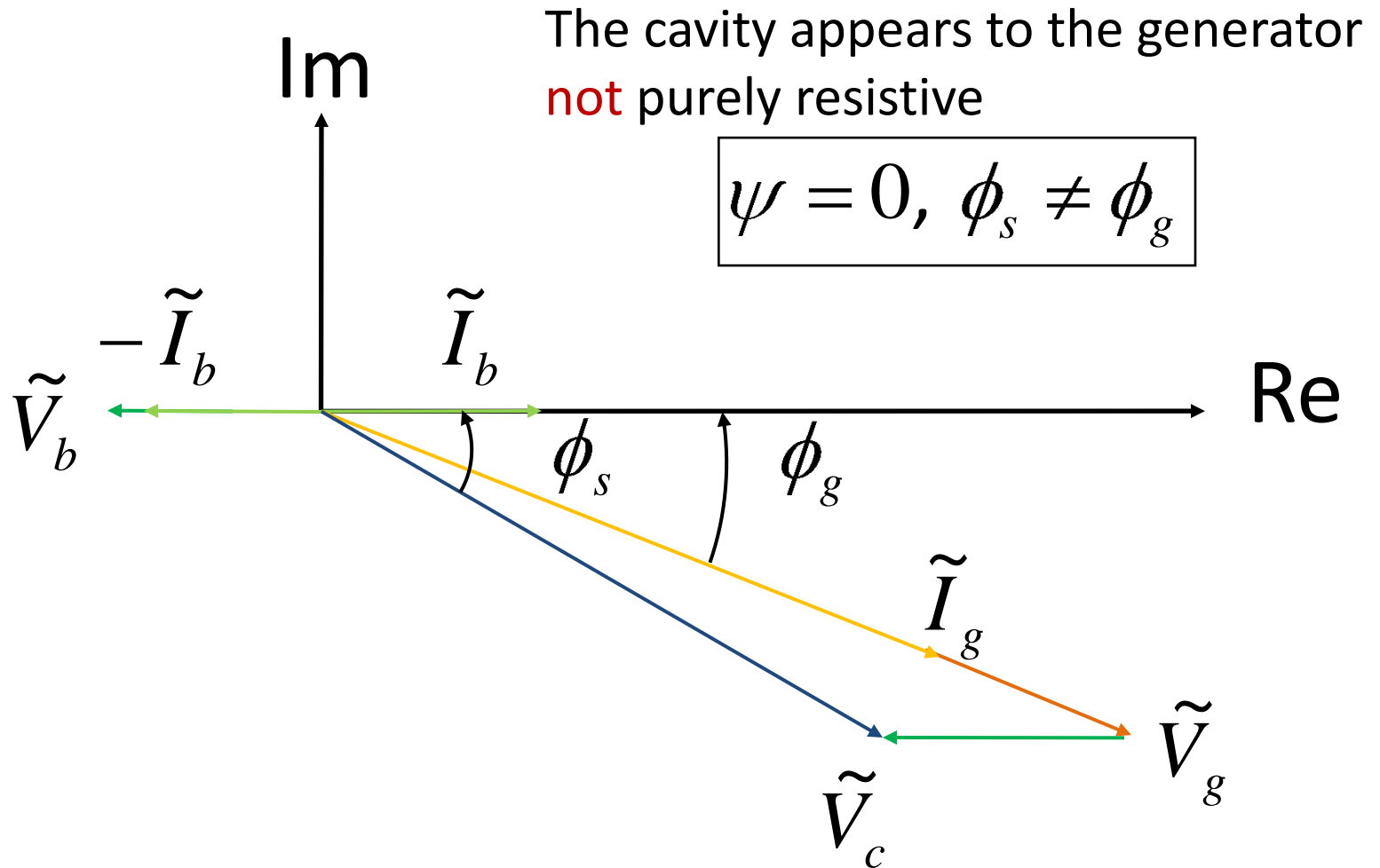
The cavity appears to the generator purely resistive

$$\psi = 0, \phi_s \approx \phi_g$$





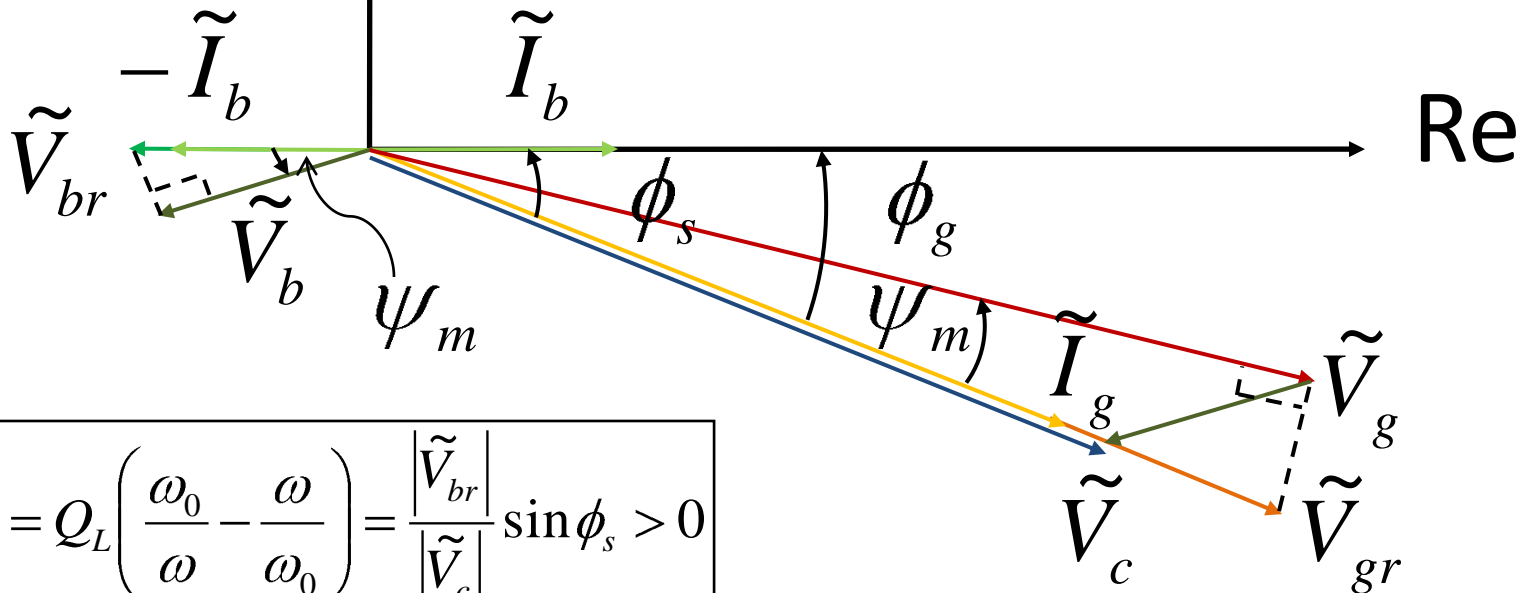
# If $V_b$ **not** negligible...



...detuning is required

The cavity appears to the generator  
purely resistive

$$\psi = \psi_m, \phi_s = \phi_g$$



$$\tan \psi_m = Q_L \left( \frac{\omega_0}{\omega} - \frac{\omega}{\omega_0} \right) = \frac{|\tilde{V}_{br}|}{|\tilde{V}_c|} \sin \phi_s > 0$$

In fact, coupling coefficient needs to be adjusted as well to further minimize reflection

# Operation Consideration

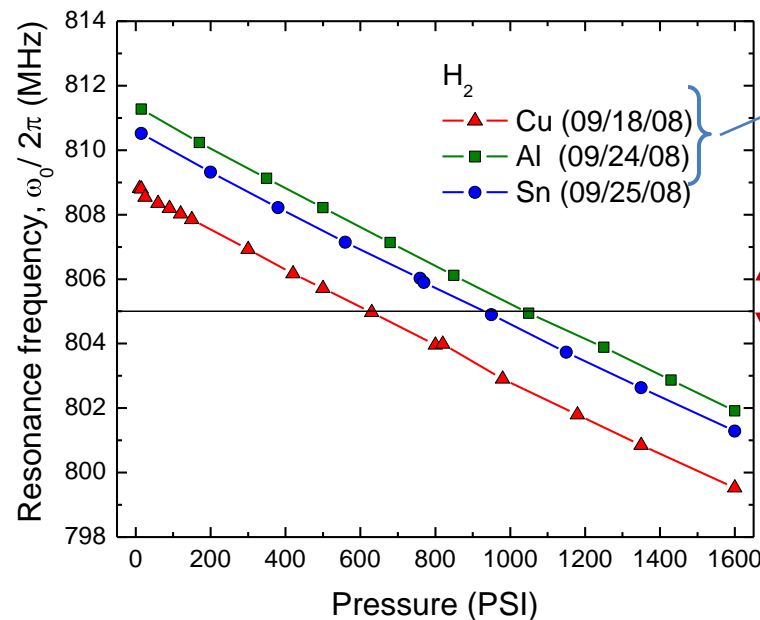
The required detuning is small

$$f_0 - f \approx \frac{f_0}{2Q_L} \tan \psi_m \sim 1.4 \text{ kHz} \ll \frac{f_0}{2Q_L} \sim 67 \text{ kHz (for 3dB)}$$

Last digit of the  
RF waveform generator

If driving frequency ( $\omega$ ) is phase-locked to the Linac for constant synchronous phase, it is important to check available pressure range.

If driving frequency ( $\omega$ ) is not phase-locked to the Linac, triggering and delay should be set properly so the RF is on when beam passes through the cavity



Shift not from material  
properties  
but from the geometry of the  
specific electrode

$f = \omega / 2\pi = 805 \text{ MHz}$   
(Driving frequency)

Install remote  
tuner next time?

## 2. Beam-induced Electrons

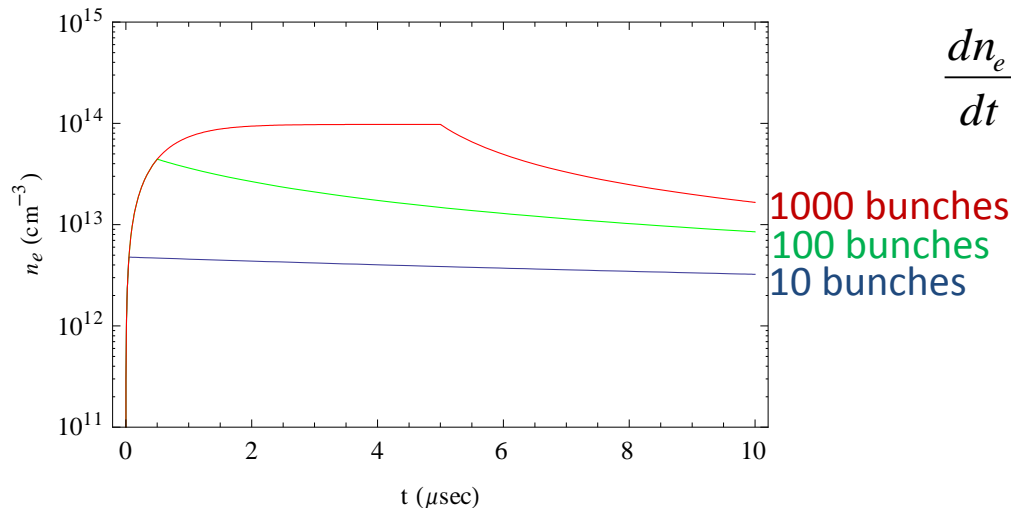
# Electron Generation

Proton beam produces electrons

by **proton-impact ionization** + **ionization by secondary electrons**

$$\frac{\Delta n_e}{1 \text{ proton}} \approx \frac{\rho(dE/dx)\Delta s}{W_i(\approx 35 \text{ eV})} \times \frac{1}{(\pi r_b^2 \Delta s)}$$

Most (or some) electrons will be thermalized quickly (energy equilibration is much faster than density evolution) by elastic + inelastic collisions, and drifting with RF until annihilated (recombination + attachment + diffusion)



$$\frac{dn_e}{dt} \approx S - \beta_r(T_e)n_e^2 - \underbrace{k_a(T_e, T_g)n_en_g - \frac{D(T_e)}{\Lambda^2}n_e}_{\text{Estimated to be small}}$$

Estimated to be small

$T_e < 1 \text{ eV}$ ,  $T_g \sim 300 \text{ K}$

→ Need more simulation

→ Oopic, LSP(Voss)

# Perturbation from Electrons

$$\sigma_{DC} \propto \frac{n_e}{\nu_m} \propto \frac{\rho}{p}$$

Need nonlinear correction for low  $E_0/p$   
 For  $E_0/p \sim 1.0$ ,  $\chi \sim 3$   
 For  $E_0/p \sim 3.9$ ,  $\chi \sim 1.5$

$$\Delta\left(\frac{1}{Q}\right) = \frac{\int_V \frac{1}{2} \sigma_{DC}(r) E_0^2(r) dV}{\omega_0 \int_V \frac{1}{2} \epsilon_0 E_0^2(r) dV} \rightarrow \chi(E_0/p) \Delta\left(\frac{1}{Q}\right)$$

Independent of  $p$

$$\Delta f = \frac{f_0}{2} \left( \frac{\omega_0}{\nu_m} \right) \times \Delta\left(\frac{1}{Q}\right) \ll 1 \text{ kHz}$$

Purely resistive:  
 no significant changes in  
 reactive component  
 ( $U_e = U_m$ )

Damping  
 coefficient

$$\gamma = Q_L \Delta\left(\frac{1}{Q}\right) = \frac{P_{Ohmic}}{P_{ext} + P_{wall}}$$

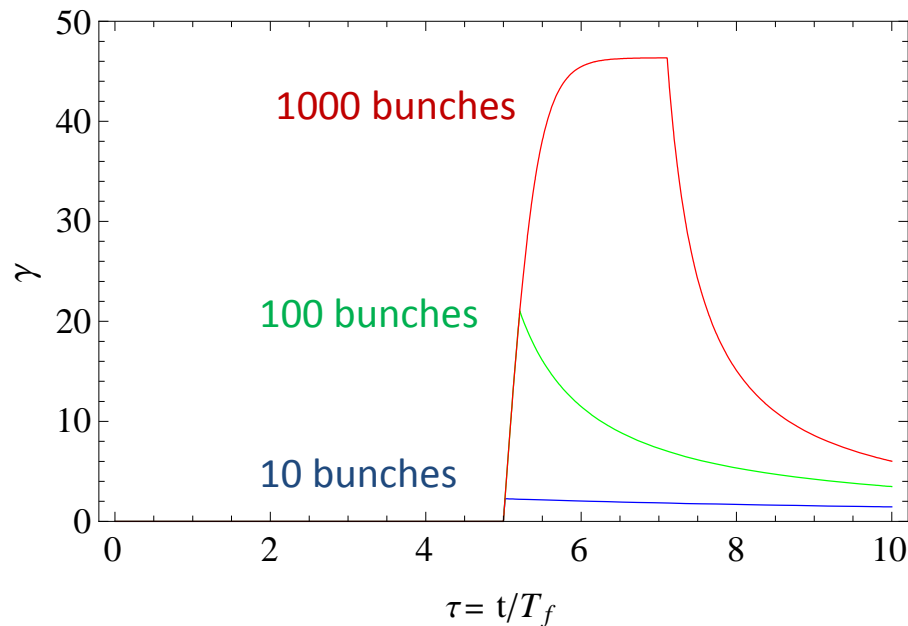
Ohmic heating of  
 hydrogen gas

# Model Equation

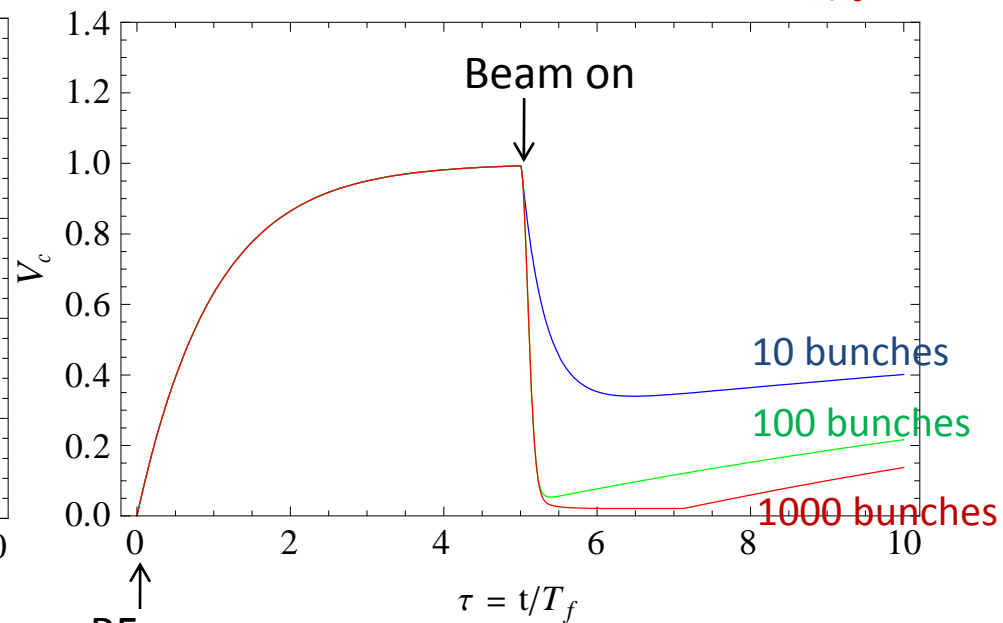
After neglecting beam current, no detuning ( $\psi = 0$ ),  
and proper normalization (for  $\gamma \rightarrow 0, \hat{V}_c(\infty) \rightarrow 1$ )

$$\frac{d\hat{V}_c}{d\tau} + [1 + \gamma(\tau)]\hat{V}_c = 1$$

Lots of reflection ( $V_R = V_C - 1$ )  
Heavily under-coupled ( $\beta_c \ll 1$ )

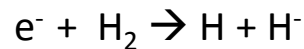


$\gamma \uparrow, T_g \uparrow, k_a \uparrow, n_e \downarrow$

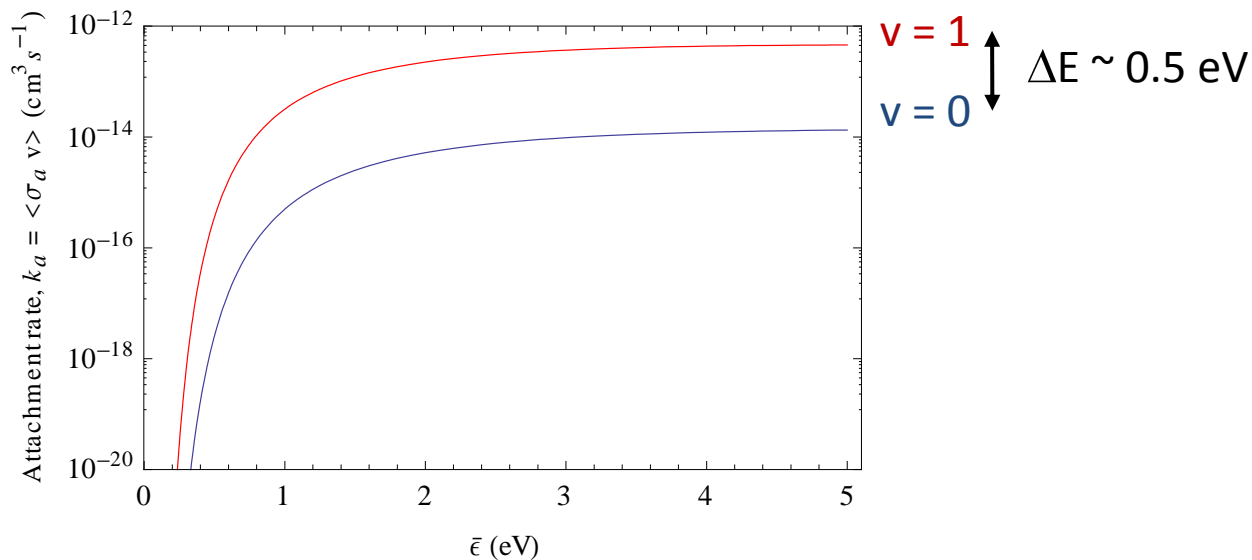


$V_c \downarrow, T_e \downarrow, \beta_r \uparrow, n_e \downarrow$

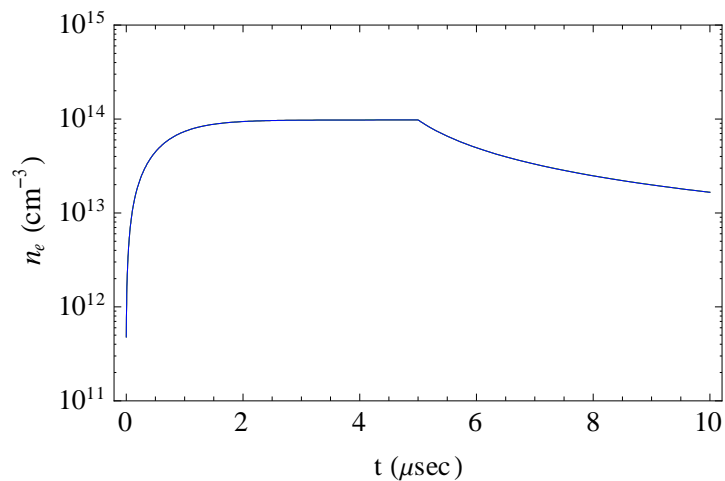
# Any Hope ?



Rate of Dissociative  
Attachment (DA)  
increases with the population  
of vibrationally  
excited hydrogen ( $v > 0$ )

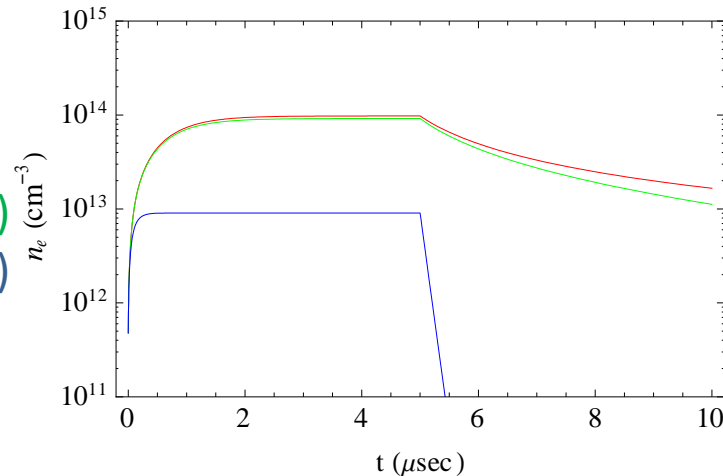


$$E_0 / p = 1.0 \text{ V/cm/torr}, \bar{\epsilon} = 0.28 \text{ eV}$$



without DA  
with DA ( $v=0$ )  
with DA ( $v=1$ )

$$E_0 / p = 3.9 \text{ V/cm/torr}, \bar{\epsilon} = 0.73 \text{ eV}$$

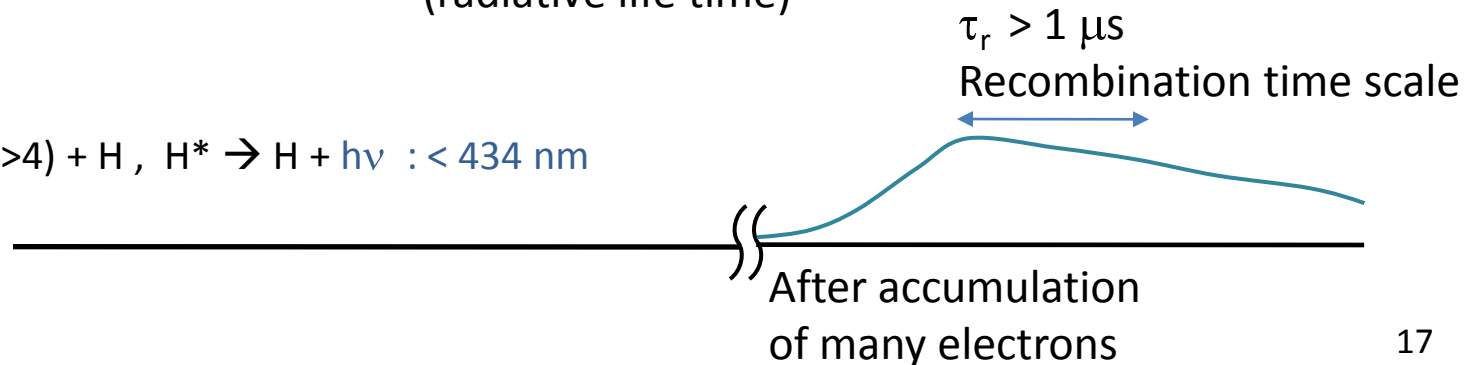
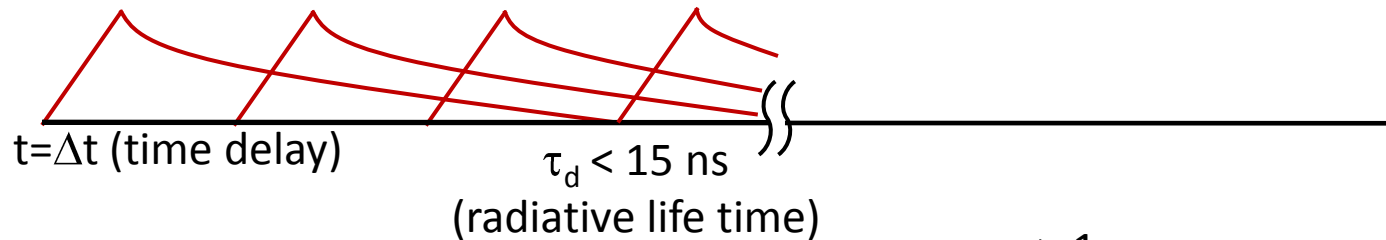
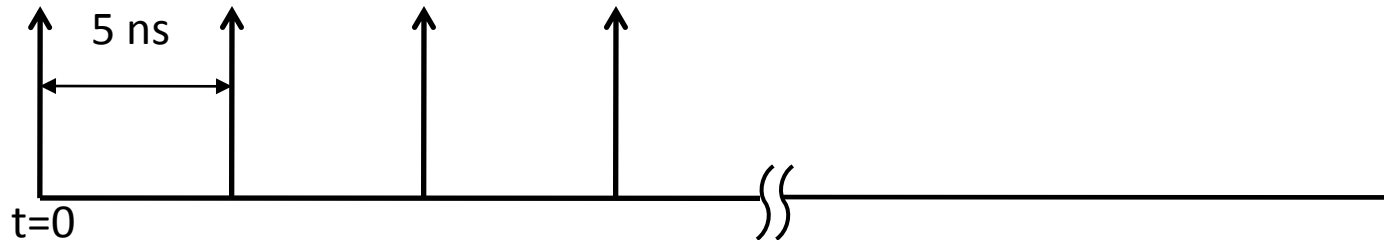




# Time Resolved Optical Emission Spectroscopy

Only if we have enough photons  $\rightarrow$  Photon counting ?

Beam pulse



# Conclusions

1. Beam loading from beam itself is **estimated to be** negligible

2. Beam loading from beam-induced electrons **is estimated to be** non-negligible

- Hopefully, recombination, attachment, and diffusion processes are significant than expected
- Hopefully, adding a small fraction of a dopant gas can help

3. Things in progress

- Beam commissioning: **Matter of time?** or **Need more manpower?**
- Simulations: **Oopic?** or **LSP?**
- Optical Measurement: **Prizm (Grating)?** or **Filter?**
- Data acquisition (< 1 shot/min):
  - ✓ BPM (or Toroid) → Beam arrival
  - ✓ LabVIEW → BPM (or Toroid), Envelope of Ref and Pickup
  - ✓ Fast scope → Ref, Pickup, **PMT1(red)**, **PMT2(blue)**

Trigger  
@ beam on

